

SOFT HIGH ENERGY SCATTERING IN NONPERTURBATIVE QCD

E. R. Berger

*LAPTH, Universite Paris-Sud, Batiment 211,
F-91405 Orsay Cedex, France
E-mail: berger@th.u-psud.fr*

In this report diffractive high energy reactions are discussed in a functional integral approach where hadronic amplitudes are calculated from vacuum expectation values of lightlike Wegner-Wilson loops. In the first part we calculate elastic differential cross sections for high energy and small momentum transfer elastic proton-proton (pp) scattering which are in reasonable agreement with the experimental data. In the second part we consider exclusive π^0 production in ep-scattering. At high energies photon and odderon exchange contribute to this reaction. We show that odderon exchange leads to a much larger inelastic than elastic π^0 production cross section, dominating the γ contribution by orders of magnitude. Observing our process at HERA would establish the soft odderon.

1 Elastic pp scattering

We will consider here some results of¹ for elastic differential cross sections $d\sigma/dt$ in elastic pp(p \bar{p}) and of² for exclusive π^0 production in ep scattering, both at high center of mass (cm) energies $\sqrt{s} \gtrsim 20$ GeV and low momentum transfer squared $|t| \lesssim O(1 \text{ GeV}^2)$ calculated within a functional approach^{3,4}.

We start with elastic pp scattering. For large energies the amplitude $T_{pp}(T_{p\bar{p}})$ is known to be dominated by $C = P = +1$ (pomeron \mathcal{P}) exchange, $C = P = -1$ (odderon \mathcal{O}) exchange^a is small or absent. In our approach, where the p is treated as a wave packet of its constituent quarks, this is achieved by considering the proton to be a quark-diquark (q-qq) system⁶, what we will do in the following. Then the result for T_{pp} is:

$$T_{pp} = (-2is) \int d^2 b_T \exp(i\vec{q}_T \cdot \vec{b}_T) \cdot \int d^2 x_T d^2 y_T |\Psi^p(x_T)|^2 |\Psi^p(y_T)|^2 \tilde{J}(\vec{x}_T, \vec{y}_T, \vec{b}_T); \quad (1)$$

The scattering amplitude (1) is obtained by first considering the scattering of quarks (antiquarks) on a fixed gluon potential G . Travelling through it the quarks (q) and diquarks (qq) pick up non-abelian phase factors. To ensure

^a The soft (nonperturbative) odderon is introduced in elastic hadron-hadron scattering as the $C = P = -1$ partner of the pomeron⁵.

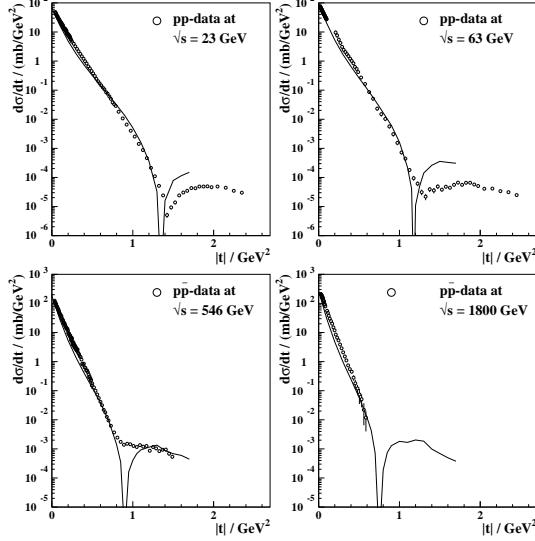


Figure 1: Differential elastic cross sections for c.m. energies $\sqrt{s} = 23, 63, 546$ and 1800 GeV.

gauge invariance the phase factors for q and qq of the same p are joined at large positive and negative times, yielding lightlike Wegner-Wilson loops W_{\pm} . The parton scattering amplitude \tilde{J} , given by the functional integral over G (indicated with brackets $\langle \rangle_G$),

$$\tilde{J}(\vec{x}_T, \vec{y}_T, \vec{b}_T) = \left\langle W_+ \left(\frac{1}{2} \vec{b}_T, \vec{x}_T \right) W_- \left(-\frac{1}{2} \vec{b}_T, \vec{y}_T \right) - 1 \right\rangle_G, \quad (2)$$

is the central object in our approach. The transverse separation between the centres of the loops (impact parameter) is given by \vec{b}_T . The vectors \vec{x}_T and \vec{y}_T give the transvers extensions and orientations of the loops. The resulting parton scattering amplitude has to be integrated over \vec{x}_T, \vec{y}_T with a measure given by the q-qq densities, $\Psi^p \sim \exp(-\vec{x}^2/(4S_p^2))$. Here S_p is the p transvers extension parameter. To calculate (1) a Fourier transform over \vec{b}_T has to be done finally.

The parton amplitude \tilde{J} is calculated in nonperturbative QCD, using the minkowski version of the model of the stochastic vacuum (MSV). In terms of a matrix cumulant expansion the result is¹:

$$\tilde{J} = \left[\frac{2}{3} \cos\left(\frac{1}{3}\chi\right) + \frac{1}{3} \cos\left(\frac{2}{3}\chi\right) - 1 \right] + i \left[-\frac{2}{3} \sin\left(\frac{1}{3}\chi\right) + \frac{1}{3} \sin\left(\frac{2}{3}\chi\right) \right]. \quad (3)$$

The real function χ is calculated from an integral over the non-local gluon condensate, for which the MSV makes an ansatz^{4,1}. The parameters that

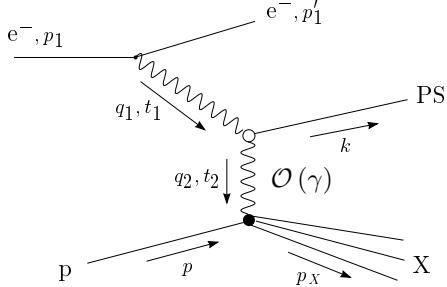


Figure 2: Feynman diagrams for pseudoscalar meson production in ep scattering at high energies with odderon (photon) exchange.

occur are the gluon condensate G_2 , the correlation lenght a and κ , a parameter measuring the non-abelian content of the condensate. In this way high energy scattering is related to the QCD vacuum parameters.

Now χ is odd under space inversion or charge conjugation, for example $\chi(-\vec{x}_T, \vec{y}_T, \vec{b}_T) = -\chi(\vec{x}_T, \vec{y}_T, \vec{b}_T)$. For this reason \tilde{J} contains \mathcal{P} and \mathcal{O} contributions. But since the q-qq densities in (1) are invariant under space inversion the \mathcal{O} contributions to T_{pp} average out consistent with experiment^b.

Leaving G_2, κ, a and S_p in (1) as free parameters we determine them in high energy scattering, using as input experimental data at $\sqrt{s} = 23$ GeV. We obtain¹: $G_2 = (529 \text{ MeV})^4, \kappa = 0.75, a = 0.32 \text{ fm}$, consistent with existing lattice measurements⁷, and a loop extension parameter S_p in the range of the electromagnetic proton radius, $S_p = 0.87 \text{ fm}$.

The term of χ proportional to κ results into string-string interaction. This causes σ_{tot} to increase with increasing extension parameter^{4,1}. Following⁴ we introduce an energy dependence of (1) by making S_p energy dependent, fixed by requiring that our model reproduces the pomeron part of σ_{tot} in the DL parametrisation (The vacuum parameters should not depend on the energy).

In Fig. 1 we show our results for $d\sigma/dt$. A major point is, that our results for $d\sigma/dt$ depend crucially on the string-string interaction, $\kappa \neq 0$, what also implies in terms of the MSV a string tension $\rho \neq 0$ and so confinement. A detailed discussion of this point can be found in¹.

2 Exclusive π^0 production in ep scattering

In this section we discuss exclusive π^0 production in high energy ep scattering. The π^0 is produced (see Fig. 2) by $\gamma\mathcal{O}$ -fusion and by $\gamma\gamma$ -fusion, but can not

^bA consequence is, that T_{pp} and $T_{\bar{p}\bar{p}}$ are equal in the diquark limit. They contain no real part which is due to a second cumulant approximation¹. As we will see, this results in infinitely deep dips in $d\sigma/dt$

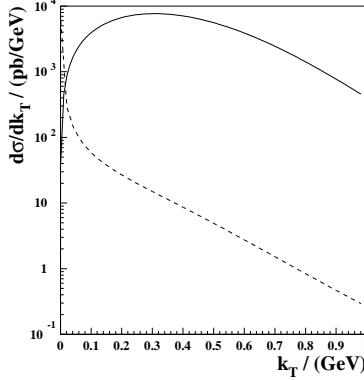


Figure 3: The k_T distribution in pion production from the 2P resonance channels for odderon exchange (solid line) compared to the complete electromagnetic result (dashed line). Interference contributions are not taken into account.

be produced by $\gamma\mathcal{P}$ -fusion since the π^0 has positive C -parity.

In Fig. 2 X stands for a proton or resonances or a sum over resonances. Here we treat the very small $Q^2 := -q_1^2$ range where the ep cross section can be calculated by folding the γp cross section with the equivalent photon spectrum of the electron.

To calculate the \mathcal{O} exchange we look at γp -scattering in the c.m. system. In the following we discuss the cases, (i) that the proton, stays intact or (ii) gets diffractively exited into the resonances $N(1520)$ with $J^P = \frac{3}{2}^-$ and $N(1535)$ with $J^P = \frac{1}{2}^-$, described as exited quark-diquark systems.

When considering unpolarised cross sections, summed over both resonances in case (ii), the quark spin degree of freedom becomes irrelevant and the calculation reduces to one where a spinless state stays intact or is exited to a 2P resonance. The amplitudes are again given by folding \tilde{J} with appropriate wave functions. For (ii) the helicity amplitudes are²

$$\begin{aligned}
 T(s_2, t_2)_{\lambda, \lambda_\gamma} &= -2is_2 \int d^2 b_T e^{iq_2 T \vec{b}_T} \hat{J}_{\lambda, \lambda_\gamma}(\vec{b}), \\
 \hat{J}(\vec{b})_{\lambda, \lambda_\gamma} &= \int \frac{d^2 x_T}{4\pi} \int \frac{d^2 y_T}{4\pi} \\
 &\times \sum_{f, h_1, h_2} \Psi_{fh_1 h_2}^* \pi^0(\vec{x}_T) \Psi_{\lambda_\gamma, fh_1 h_2}^\gamma(\vec{x}_T) \\
 &\times \Psi_\lambda^{*2P}(\vec{y}_T) \Psi^p(\vec{y}_T) \tilde{J}(\vec{x}_T, \vec{y}_T, \vec{b}_T).
 \end{aligned} \tag{4}$$

Here $\lambda_\gamma(\lambda)$ is the helicity and $\Psi^\gamma(\Psi^{2P})$ the wave function of the photon (2P state). Now, when the proton stays intact there occurs in (4) the quark-diquark density ($\Psi_\lambda^{2P} \rightarrow \Psi^p$). But, since this density is symmetric under a

parity transformation whereas the odderon coupling changes sign, there is a cancellation when we integrate over all angles. On the other hand when the proton gets excited to a negative parity state like the resonances $N(1520)$ and $N(1535)$ there is no cancellation (the overlap is odd under parity) and the odderon couples to the nucleon without any restriction^{2,8}. Thus in our model the odderon couples only if breakup of the proton occurs.

The photoproduction cross section calculated from (4) is independent of s_2 and so the EPA conversion to electroproduction can be achieved by simply multiplying it with a constant $c_{\text{EPA}} = 0.0136$.

In the following we consider only ep scattering. An experimentally preferred observable is the k_T spectrum, the transverse momentum distribution of the π^0 with respect to the beam direction. This is displayed in Fig. 3. We also show there the distribution of the pion's k_T for γ exchange summed over the elastic and all inelastic channels (with invariant mass $M_X \leq 2$ GeV). As we can see γ exchange is larger than \mathcal{O} exchange only for very small k_T . For $k_T \gtrsim 0.1$ GeV the \mathcal{O} exchange dominates by orders of magnitude.

3 Conclusions

The functional integral approach gives a reasonable description of soft elastic pp scattering at high energies, relating it to fundamental parameters of the QCD vacuum. The values determined from high energy scattering compare well with lattice results. In π^0 production for the p breakup into negative parity hadronic final states, a large odderon cross section arises which should be observable at HERA. This would establish the soft odderon as an exchange-object in high energy scattering on an equal footing with the soft pomeron.

References

1. E. R. Berger, O. Nachtmann, Eur. Phys. J. C7, 459 (1999)
2. E. R. Berger, A. Donnachie, H. G. Dosch, W. Kilian, O. Nachtmann, M. Rueter, Eur. Phys. J. C9, 491 (1999);
3. O. Nachtmann, Ann. Phys. 209 (1991) 436.
4. H. G. Dosch, E. Ferreira, A. Krämer, Phys. Rev. D50 (1994) 1992.
5. L. Lukaszuk, B. Nicolescu, Nuov. Cim. Lett. 8, 405 (1973); D. Joynson, E. Leader, C. Lopez, B. Nicolescu, Nuov. Cim. A 30, 345 (1975)
6. M. Rueter, H. G. Dosch, Phys. Lett. B 380, 177 (1996)
7. E. Meggiolaro, Phys. Lett. B451, (1999) 414
8. M. Rueter , H.G. Dosch , O. Nachtmann, Phys. Rev. D59, 014018, 1999